

Autoresonant transition in the presence of noise and self-fields

Ido Barth¹, Lazar Friedland¹, Eli Sarid², and Arkadi Shagalov³

1. Racah Institute of Physics, Hebrew University of Jerusalem, Jerusalem, Israel.

2. Department of Physics, NRCN-Nuclear Research Center Negev, Beer Sheva, Israel.

3. Institute of Metal Physics, Ekaterinburg, Russian Federation.

Abstract

Autoresonance is a nonlinear phenomenon characterized by continuing phase-locking between a dynamical system and external oscillatory perturbations despite slow variation of system's parameters. A sharp threshold for autoresonant capture of an ensemble of trapped particles driven by chirped frequency oscillations is analyzed. It is shown that at small temperatures T , the capture probability versus driving amplitude is a smoothed step function with the step location and width scaling as $\alpha^{3/4}$ (α being the chirp rate) and $(\alpha T)^{1/2}$, respectively. Strong repulsive self-fields reduce the width of the threshold considerably, as the ensemble forms a localized autoresonant macroparticle.

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Introduction

A. Nonlinear oscillator

A driven, damped, noisy pendulum (dimensionless variables).

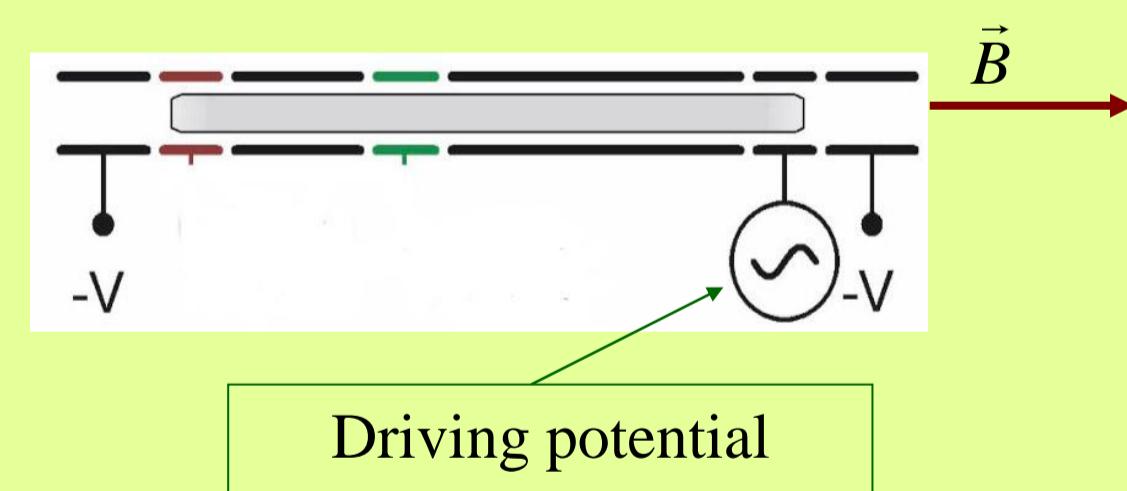
$$\ddot{x} + \nu \dot{x} + \sin x = \varepsilon \cos \psi_d + \sqrt{2\nu T} \xi(t)$$

- Delta correlated Gaussian noise: $\langle \xi(0)\xi(t) \rangle = \delta(t)$
- Chirped driving frequency: $\omega_d(t) = \dot{\psi}_d = 1 - \alpha t$ passes through linear resonance at $t = 0$.
- A continuing *phase-locking* is established and preserved due to nonlinearity.
- $T=0 \rightarrow$ sharp threshold for capture into resonance for $\varepsilon > \varepsilon^{cr}$.

What is the effect of the thermal noise?

B. Trapped nonneutral plasma

Penning trap:



Driving potential

Vlasov-Poisson equations:

$$\begin{cases} f_t + u f_x - (V_x + \varphi_x + \varphi_x^{drive}) f_u = 0 \\ \varphi_{xx} - \chi^2 \varphi = -\eta^2 \int_{-\infty}^{\infty} f(u, x, t) du \end{cases}$$

- External potential: $V = 1 - \cos x$
- Driving potential: $\varphi_d = \varepsilon x \cos \psi_d$
- Self-potential: $\varphi(x)$

What is the effect of the repulsive self-field?

Theory

Assumptions

- Sufficiently small temperature \rightarrow Weak nonlinearity.
- Sufficiently fast passage through resonance \rightarrow Noise is negligible.
- However, noise and dissipation form thermally distributed, random initial conditions, f_0 (FDT).

Analysis

- Seek solution of form $x = a \cos \theta$.
- Phase mismatch: $\phi = \theta - \psi_d$
- Neglect non resonant terms (**single resonance approximation**).
- Average over the fast phase.

$$\begin{cases} \dot{a} = -\frac{\nu}{2} a - \frac{\varepsilon}{2} \sin \phi \\ \dot{\phi} = \frac{a^2}{16} - \alpha t - \frac{\varepsilon}{2a} \cos \phi \end{cases}$$

- Rescale the time: $t \rightarrow \tau = \sqrt{\alpha} t$.

$$\text{Define } \gamma = \frac{\nu}{2\sqrt{\alpha}}, \quad A = \frac{a}{4\alpha^{1/4}}, \quad \text{and } \mu = \sqrt{\frac{3}{36}} \frac{\varepsilon}{\alpha^{3/4}}.$$

- Complex variable: $\Psi = A e^{i\phi}$

$$\text{NLSE: } i\Psi + (|\Psi|^2 - \tau)\Psi + i\gamma\Psi = \mu$$

- Phase locked asymptotic solution for $\mu > \mu_{cr}(\gamma, A_0, \phi_0)$:

$$|\Psi|^2 \approx \tau, \quad \phi = \text{const}$$

- Critical driving parameter is periodic in the initial phase mismatch

$$\mu_{cr} = c_0(\gamma) + \kappa A_0 \cos \phi_0$$

$$\text{Resonant capture probability: } P(\mu, A_0) = \begin{cases} 0 & \mu < c_0 - \kappa A_0 \\ \frac{1}{\pi} \arccos\left(\frac{c_0 - \mu}{\kappa A_0}\right) & |\mu - c_0| < \kappa A_0 \\ 1 & \mu > c_0 + \kappa A_0 \end{cases}$$

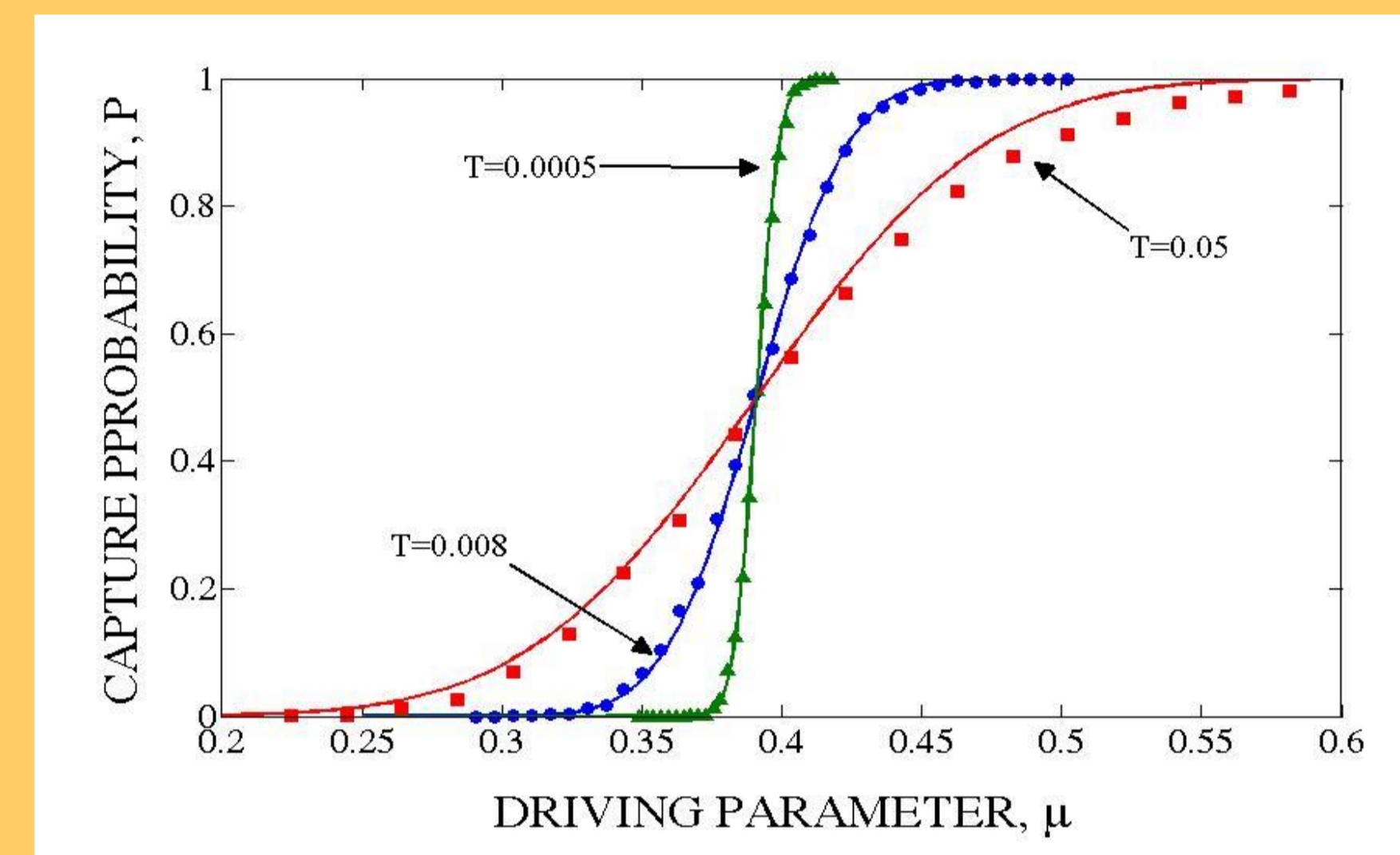
- Integration over the thermal distribution

$$P(\mu) = \int_0^{\infty} P(\mu, A_0) f_0(A_0) dA_0$$

Results

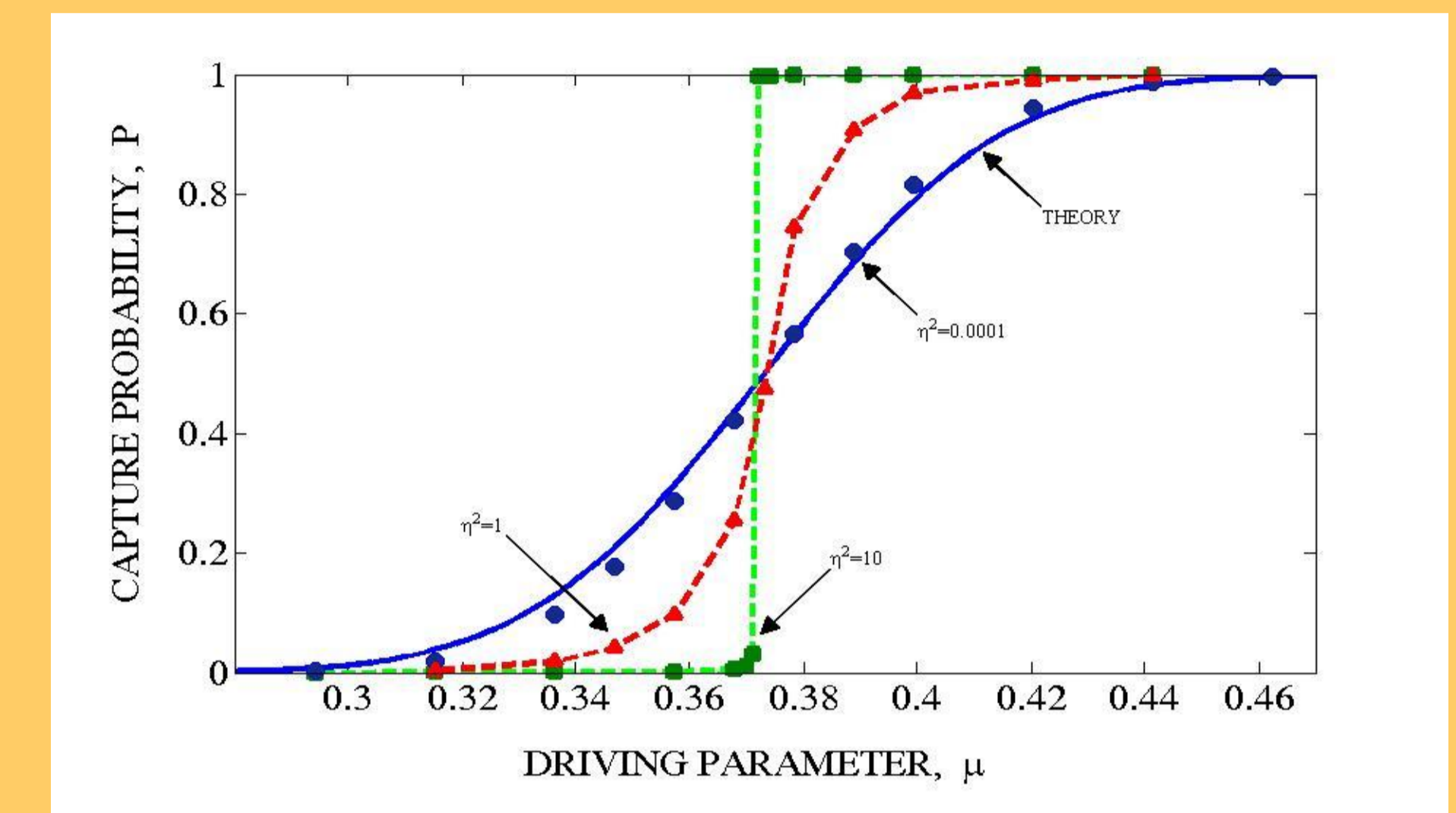
Probabilities of capture into autoresonance

A. Nonlinear oscillator



Transition width increases with temperature !

B. Trapped plasma

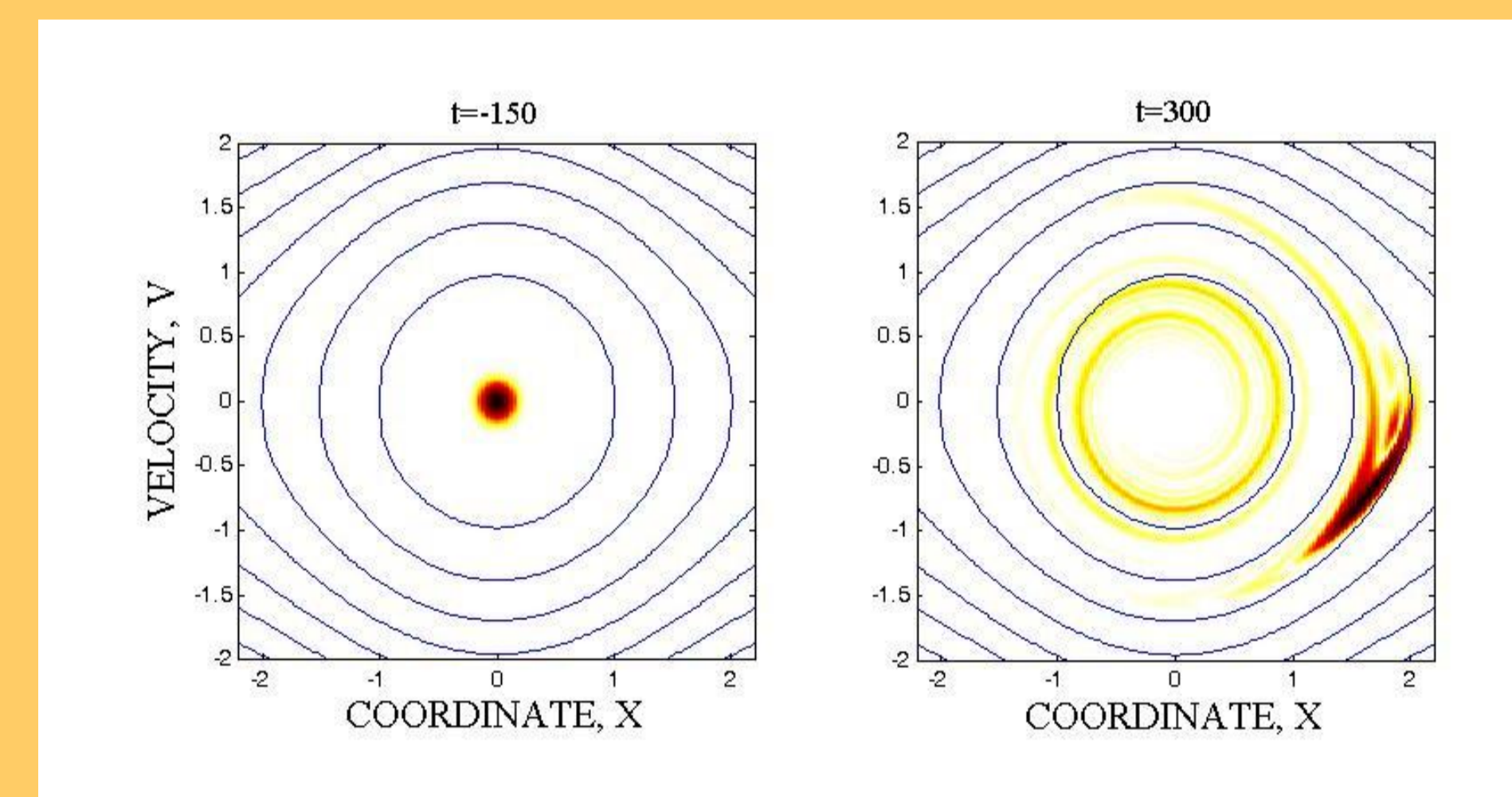


Repulsive self-field makes the width narrower !

Plasma distribution in phase space

Low plasma density ($\eta^2 = 0.0001$)

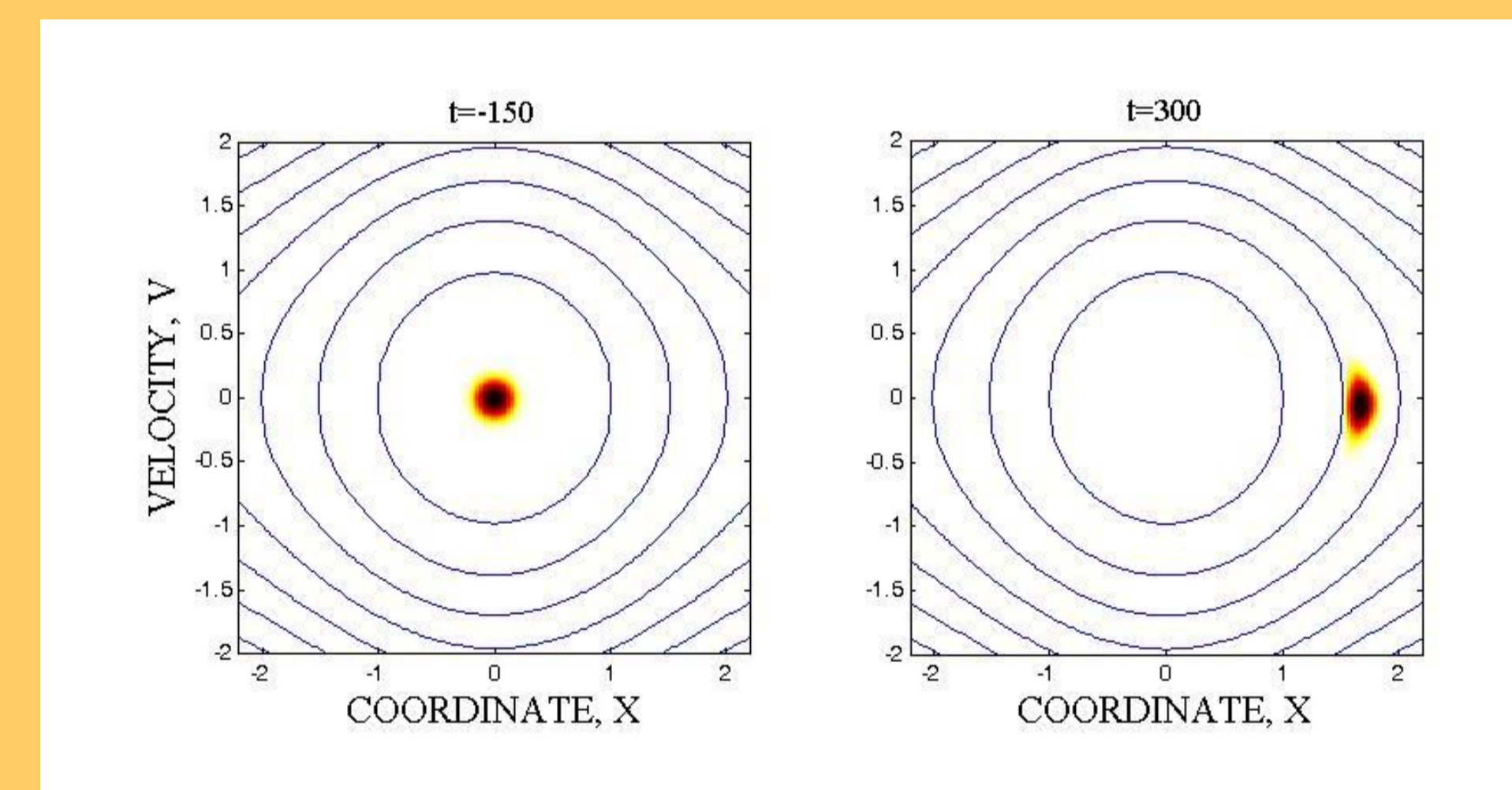
$\mu = 0.38$



57 % captured

High plasma density ($\eta^2 = 10$)

$\mu = 0.38$



100 % captured

Autoresonant macroparticle

Conclusions

- We presented a theory describing the shape and the width of the autoresonant phase-locking transition versus the driving amplitude for both stochastic oscillators and dilute nonneutral plasmas in Penning traps driven by a chirped frequency drive.
- We observed and explained why the repulsive self-field yields a strong bunching effect and a significant narrowing of the phase-locking transition width.
- The plasma in this case behaves like a single autoresonant macroparticle and can be efficiently controlled by external parameters.
- Our theory agrees with simulations for the driven oscillator in the presence of thermal noise, as well as with Vlasov-Poisson simulations in the case of sufficiently low plasma densities.

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